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Ricci Flow

$$\underbrace{(g_{ij})_t}_{\text{Riemannian tensor}} = -2 \underbrace{R_{ij}}_{\text{Ricci tensor}}$$

1982 Richard Hamilton - short time existence (smooth)

later: $n=3$, $R_{ij} > 0$ initially \Rightarrow Solution blows up in finite time
and \Downarrow

Renormalized Ricci flow (constant volume = 1) \Rightarrow Solution converges exponentially fast to a metric of constant curvature
 $S^3/\pi \leftarrow M^3$

"Technique which led to this remarkable result":

Hamilton observed Ricci flow is analagous to the heat equation (parabolic)

Look at a length minimizing geodesic joining two pts $\gamma: (0, L) \rightarrow M^3$



Study variation $\frac{d}{ds} \int_0^L \exp \rho(t) (s, \gamma(t)) ds$

Essentially like a heat eqn for the length of short geodesics.

Evolution Equ for the curvature (Hamilton) 2

$$\frac{\partial}{\partial t}(R \dots) = \Delta_g R \dots + Q \dots$$

curvature
tensor
← R...

(also parabolic)

Curvature is getting more positive
along the flow

scalar curvature R satisfies

~~Equation~~
$$\frac{\partial}{\partial t} R = \Delta R + 2(R_{ij})^2$$

Warning $n > 3 \rightarrow$ Positive Ricci is not preserved
only the positivity of the curvature
operator

Afterwards

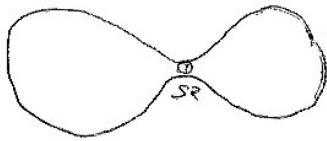
Hamilton + Others tried to extend
Ricci flow technique without
assumptions of positivity
of Ricci curvature.

In general

Flow doesn't exist \forall time even renormalized except in dimension 2 (Hamilton).

Einstein Metrics are limits when?

(note in $\dim=3$, this implies limit has const sect curv)



pinching
 \Rightarrow Singularity develops (in a short time)



How to continue through the singularity?

Hamilton 1996, $n=4$, positivity on certain curvatures,



remove singularity



restart flow

(Does this only deal w/ pt singularities?)

Ultimately hope flow stops & reassemble initial manifold using connected sums.

Naduehy

Want: $n=3$ arbitrary (M^3, g_0) via Ricci flow,

cut & paste to remove singularity,

& get M^3 as connected sums of manifolds of constant sectional curvature

Topologists: Geometric decomposition using incompressible tori. (not connected sums?)

Hamilton 1999:

$n=3$, analyze Ricci flow if it exists \forall time with bounded normalized curvatures \Rightarrow Geometrization of such M^3

* ————— *

Now a new technique which works in all

dimensions - Perelman (complements Hamilton's Max Principle (inspired by it))

Integral Monotonicity Formulas

(Ptwise!)

Warning: (in response to a question) M^3 has Ricci \geq below then M^3 has some topological finiteness

In general

Does a parabolic flow come from a gradient flow or not?

Ricci flow does in a certain sense!

$$F(g_{ij}, f) = \int_M (R + |\nabla f|^2) e^{-f} dV$$

$f: M \rightarrow \mathbb{R}$
 g_{ij} metric on M

R scalar curvature

dV Riem volume

Look at the variational formula for F .

$$S\mathcal{F}(v_{ij}, h) = \int_M [-v_{ij} (R_{ij} + \nabla_i \nabla_j f) + (\frac{v}{2} - h) (2\Delta f - |\nabla f|^2 + R)] e^{-f} dv$$

$$v_{ij} = Sg_{ij} \quad h = Sf \quad v = g^{ij} v_{ij}$$

Consider variations where $\frac{v}{2} - h = 0$.

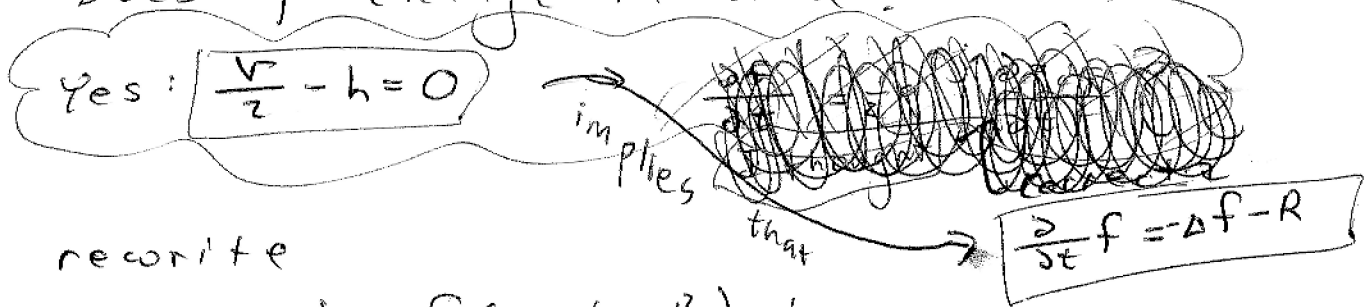
$$S\mathcal{F}(v_{ij}, h) = \int_M [-v_{ij} (R_{ij} + \nabla_i \nabla_j f) (2\Delta f - |\nabla f|^2 + R)] e^{-f} dv$$

So the variational flow in this direction

$$\frac{d}{dt} g_{ij} = -2 (R_{ij} + \nabla_i \nabla_j f)$$

Differs from Ricci flow by a diffeomorphism

Does f change in time?



Let's rewrite

$$F(g_{ij}) = \int_M (R + |\nabla f|^2) dm$$

fixed measure

~~define f so that~~ $dm = e^{-f} dv$

volume form of g

Fix a measure, find an f , compute the flow
 Fix a different measure \rightarrow still a diffeomorphic flow.
 (Like gauge theory)

Note $F(g_{ij}) = \int_M (R + |\nabla f|^2) dm$
 where $f : dm = e^{-f} dv$

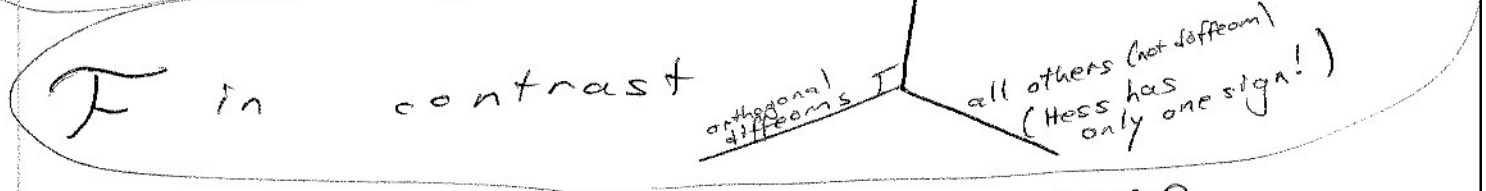
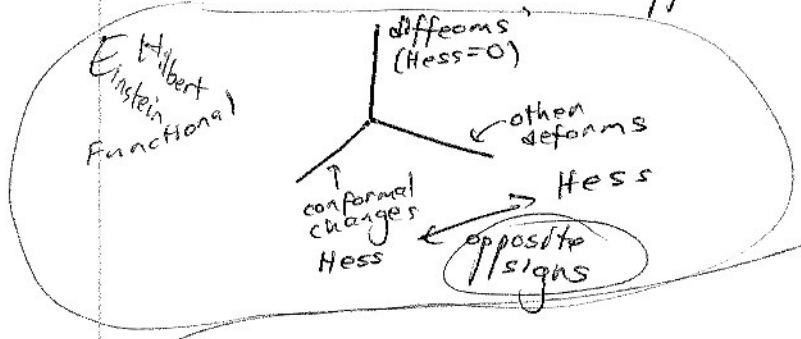
(f is a function of g_{ij} ptwise $-\log(\det g_{ij})$)

Warning doesn't necessarily exist even for a short time due to $f_t = -\Delta f - R$ difficulties. (like a backwards heat eqn!)

Consider Einstein Hilbert functional

$$E.H = \int_M R dv$$

Gradient flow is $\frac{\partial}{\partial t} |g_{ij}| = Rg_{ij} - 2R_{ij}$
 this doesn't exist for a short time for most g_{ij} because it is diffeom invariant,



F is only invariant wrt measure preserving diffeomorphisms.



~~Once we have~~

Once you have a gradient flow you can get a monotonicity formula

$$\frac{\partial}{\partial t} (g_{ij}) = -2(R_{ij} + \nabla_i \nabla_j f) \Rightarrow \frac{\partial}{\partial t} \mathcal{F} \geq 0$$

$$\frac{\partial}{\partial t} f = -\Delta f - R$$

But this flow may not exist for a short time, so look at Ricci flow with an appropriate f

$$\left. \begin{aligned} \frac{\partial}{\partial t} (g_{ij}) &= -2R_{ij} \\ f_t &= -\Delta f + |\nabla f|^2 - R \end{aligned} \right\} \frac{\partial}{\partial t} \mathcal{F} \geq 0$$

Change of variables $u = e^{-f}$

$$\frac{\partial u}{\partial t} = -\Delta u + Ru$$

"Adjoint Heat Eqn"

If $\frac{\partial h}{\partial t} = \Delta h$ then

$$\int_M h u^{\text{dV}} = \text{const.}$$

Now get rid of f (or u) in the monotonicity formula by minimizing over f :

$$\text{Let } \lambda(g_{ij}) = \inf_f \mathcal{F}(g_{ij}, f) \text{ s.t. } \int e^{-f} \text{dV} = 1$$

then under Ricci flow $\frac{d}{dt} \lambda \geq 0$ (Remarks? inf exists by heat page \rightarrow)

~~Let $\lambda(g_{ij}) = \inf_f \mathcal{F}(g_{ij}, f)$ s.t. $\int e^{-f} \text{dV} = 1$~~

$$\lambda(g_{ij}) = \inf_{\substack{f \in C^\infty \\ \int_M \omega^2 dV = 1}} \mathcal{F}(g_{ij}, f)$$

$\omega = e^{-\frac{1}{2}f}$
substitute

$$= \inf_{\substack{\omega \\ \int_M \omega^2 dV = 1}} \int_M (R\omega^2 + 4|\nabla\omega|^2) dV = \inf_{\omega} \frac{\int_M R\omega^2 + 4|\nabla\omega|^2 dV}{\int_M \omega^2 dV}$$

= lowest eigenvalue of $-4\Delta + R$
(Rayleigh Ritz Quotient). ☺

So $\frac{d}{dt} \lambda(g_{ij}) \geq 0$ where $\lambda =$ lowest eigenvalue of $-4\Delta + R$.

Note λ is not a scale invariant quantity!
 This allows Ricci flow to have the freedom to return to a starting metric that has just been rescaled (homothety). (Sort of like being periodic without being a soliton)

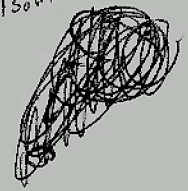
So we try to study something scale invariant.
 $\bar{\lambda} = \lambda V^{2/n}$ is scale invariant.

In fact $\frac{d}{dt} \bar{\lambda} \geq 0$ if $\bar{\lambda} \leq 0$

So if g has $\bar{\lambda} < 0$ then we cannot come back to ourselves under homothety. (Except ∇^2 soliton case?)

(Does this correspond to $R < 0$ case?)

$R = 0 \Rightarrow \bar{\lambda} > 0$ since -4Δ is a nonnegative operator.



Now study

$$W(g_{ij}, f, \tau) = \int_M (4\pi\tau)^{-\frac{n}{2}} e^{-f} dV [\tau(R + |\nabla f|^2) + f - \frac{1}{2}R]$$

scale invariant wrt $g + \tau$ (together?)

$$\frac{d}{dt} g_{ij} = -2R_{ij}$$

$$u_t = -\Delta u + Ru \quad u = (4\pi\tau)^{-n/2} e^{-f}$$

$$\tau_t = -1$$

Then $\frac{d}{dt} W \geq 0$

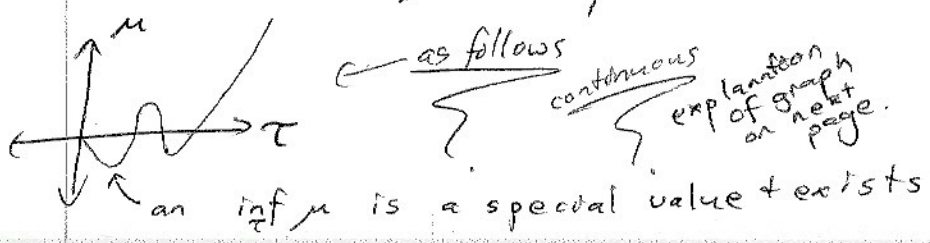
$$\mu(g_{ij}, \tau) = \inf_{\int u=1} W(g_{ij}, \tau) *$$

$$\frac{d}{dt} g_{ij} = -2R_{ij} + \tau_t = -1 \implies \frac{d}{dt} \mu \geq 0$$

Is μ some kind of eigenvalue?

Yes, for a nonlinear operator, ~~so~~ so it is not easy to prove that $\inf *$ exists. But can do this.

Trouble is μ depends on the scale parameter, τ .



Rules out homothety issue on page 8.

$\frac{d}{dt} \inf_{\tau} \mu \geq 0$

Let $g_{ij}(0) = \bar{g}_{ij}$ + run Ricci flow on $[0, T]$.

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Let $u|_{t=T} = \delta$ delta function + $\tau = 0$ at $t = T$

Solve $u_t = -\Delta u + Ru$ ← backwards? Yes in time on $(0, T)$
 Calculate W + get $W \leq 0$

So μ has to be < 0 on this time interval. $[0, T]$

The behavior of μ at $\tau = 0$ is equivalent to computing $\frac{u}{\text{Euclidean metric of } \tau = \frac{1}{2}}$ why? ?

minimize:

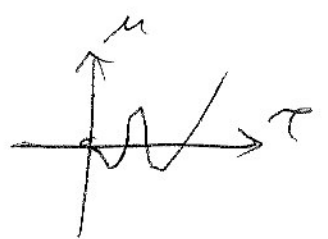
$$\min_{\int_{\mathbb{R}^n} (2\pi)^{-n/2} e^{-f} = 1} \int_{\mathbb{R}^n} (2\pi)^{-n/2} e^{-f} \left(\frac{1}{2} |\nabla f|^2 + f - n \right)$$

Shown by the Logarithmic Sobolev Inequality

$$d\mu = \text{Gaussian measure on } \mathbb{R}^n = (2\pi)^{-n/2} e^{-\frac{1}{2}|x|^2} dx_1 \dots dx_n$$

$$\int |\nabla \phi|^2 d\mu \geq \frac{1}{4} \int \phi^2 \log \phi^2 d\mu$$

minimum is 0 when $f = \frac{1}{2}|x|^2$ something?

So now we have justified 
 (at least behavior near $\tau = 0$)

In Statistical Mechanics

$$\tau, \beta = \tau^{-1}, \quad Z = \int_{\mathbb{R}} \exp(-\beta E) d\omega(E)$$

distribution over \mathbb{R} of values of energy called density of states.

$$\langle E \rangle = - \frac{\partial}{\partial \beta} \log(Z)$$

↑
average energy

$$\int_{\mathbb{R}} E \exp(-\beta E) d\omega(E)$$

Entropy $S = \beta \langle E \rangle + \log Z$

$$= \sum_i -p_i \log p_i$$

$p_i = \text{states}$

Dispersion $\sigma^2 = \langle (E - \langle E \rangle)^2 \rangle$

$$= \frac{\partial^2}{(\partial \beta)^2} \log Z$$

Now relate Ricci flow to this picture

$$(g_{ij})_{\tau} = 2(R_{ij} + \nabla_i \nabla_j f)$$

$$dm = u dv$$

$$u = (4\pi\tau)^{-n/2} e^{-f}$$

$$\log Z = \int (-f + \frac{n}{2}) dm$$

Is this compatible with the Statistical Mechanics?
I make no such claim. This is just
a formal analogy - Perelman

Let $\beta = \tau^{-1}$

$$\langle E \rangle = \tau^2 \int_M (R + |\nabla f|^2 - \frac{n}{2\tau}) dm$$

$$S = - \int_M (\tau(R + |\nabla f|^2) + f - n) dm = -W(g_{ij}, f, \tau)$$

$$\sigma = 2\tau^4 \int_M (R_{ij} + \nabla_i \nabla_j f - \frac{1}{2\tau} g_{ij})^2 dm$$

$\sigma > 0 \iff$ monotonicity of S ?

If $u = \delta$ function at $T=0$ then $\langle E \rangle$ is positive.

I think
This is 0
iff you have
a gradient
shrinking Ricci
soliton

$$\frac{d}{dt} F \geq \frac{2}{n} F^2$$

Question: What can you say about a system with bounded entropy?

In dimension 2, get nice properties for σ which don't hold in dimension 3.

Now some physics stipulations:

$$\mathcal{F}(g_{ij}, f) = \int_M (R + 10|f|^2) e^{-f} dV$$

↑
Dilaton? Field

This functional is used in String Theory

Physicists usually add an antisymmetric two form but so far it seems unrelated to Ricci flow.

Gradient + flow $(g_{ij})_t = 2(R_{ij} + \nabla_i \nabla_j f)$

Renormalization ~~Group~~ ~~Flow~~ ~~Flow~~

Ricci flow = 1st approx for 2 diml nonlinear σ model Friedman, Annals of Physics 1985

Metaphysical level:

Regarding physics at many scales with different coupling constants for each scale. Hierarchy of scales

(ultraviolet \longleftrightarrow infrared)

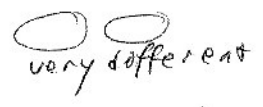
distance small
energy large

\longleftrightarrow
uncertainty principle

Calculate at a higher energy scale & then average to find out behavior at a lower scale. W is this expectation. Z suppresses contribution of higher energies.

Paradox:

Higher Energy (Strong Microscope)



Global Level

(I'm not sure what he's getting at here.)

* ————— *

Wednesday
Monotonicity
Applications ~ Bishop Gromov
→ Geometric Results.
(no local collapse).